

COMMUTING FAMILIES OF SYMMETRIC DIFFERENTIAL OPERATORS

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INTRODUCTION

Many commuting families of differential operators or completely integrable quantum systems have been constructed in connection with root systems (cf. [OP2] and references therein). Such families often have a certain symmetry in coordinates.

The radial parts of invariant differential operators on symmetric spaces give a good example of a commuting family of differential operators (cf. [HC]). In this case some parameters take only some discrete values determined by the dimensions of the root spaces for the symmetric spaces. On the other hand, [Sj] generalized them to complex parameters for the root system of type A_n . The same generalization was given by [H1], [H2], [HO], [Op1], [Op2] for general root systems. If the root system is of classical type, their operators give examples of the commuting families studied in this note (cf. Remark 3 iii). Namely we shall determine all the families under the assumption of a symmetry in coordinates.

Let W be the Weyl group of type A_{n-1} with $n \geq 3$ or of type B_n with $n \geq 2$ or of type D_n with $n \geq 4$. We identify W with the group of the coordinate transformations

$$(x_1, \dots, x_n) \mapsto (\varepsilon_1 x_{\sigma(1)}, \dots, \varepsilon_n x_{\sigma(n)})$$

of \mathbb{R}^n , where σ are the elements of the n -th symmetric group \mathfrak{S}_n and

$$\begin{cases} \varepsilon_1 = \dots = \varepsilon_n = 1 & \text{if } W \text{ is of type } A_{n-1}, \\ \varepsilon_1 = \pm 1, \dots, \varepsilon_n = \pm 1 & \text{if } W \text{ is of type } B_n, \\ \varepsilon_1 = \pm 1, \dots, \varepsilon_n = \pm 1 \text{ and } \#\{i; \varepsilon_i = -1\} \text{ is even} & \text{if } W \text{ is of type } D_n. \end{cases}$$

We examine the Laplacian

$$P = -\frac{1}{2} \sum_{1 \leq j \leq n} \frac{\partial^2}{\partial x_j^2} + V(x)$$

on \mathbb{R}^n with a W -invariant potential $V(x)$ which has enough W -invariant commuting differential operators. To be precise we assume that there exist W -invariant differential operators P_1, \dots, P_n with

$$P \in \mathbb{C}[P_1, \dots, P_n]$$

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and

$$[P_i, P_j] = 0 \quad \text{for } 1 \leq i < j \leq n$$

such that

$$P_j = \sum_{1 \leq i_1 < \dots < i_j \leq n} \partial_{i_1} \cdots \partial_{i_j} + R_j \quad \text{with } \text{ord } R_j < j \quad \text{for } 1 \leq j \leq n$$

or

$$P_j = \sum_{1 \leq i_1 < \dots < i_j \leq n} \partial_{i_1}^2 \cdots \partial_{i_j}^2 + R_j \quad \text{with } \text{ord } R_j < 2j \quad \text{for } 1 \leq j \leq n$$

or

$$\begin{cases} P_n = \partial_1 \cdots \partial_n + R_n & \text{with } \text{ord } R_n < n, \\ P_j = \sum_{1 \leq i_1 < \dots < i_j \leq n} \partial_{i_1}^2 \cdots \partial_{i_j}^2 + R_j & \text{with } \text{ord } R_j < 2j \quad \text{for } 1 \leq j \leq n-1, \end{cases}$$

if the type of W is A_{n-1} or B_n or D_n , respectively. Here $\mathbb{C}[P_1, \dots, P_n]$ is the commutative algebra over \mathbb{C} generated by P_1, \dots, P_n and for simplicity we put $\partial_i = \frac{\partial}{\partial x_i}$ and $\text{ord } R_j$ are the orders of differential operators R_j .

In this note, we assume that the coefficients of the differential operators are extended to holomorphic functions on a Zariski open subset Ω' of an open connected neighborhood Ω of the origin of the complexification \mathbb{C}^n of \mathbb{R}^n . Namely there exists a non-zero holomorphic function ϕ on Ω with $\Omega' = \{x \in \Omega; \phi(x) \neq 0\}$.

DETERMINATION OF THE COMMUTING FAMILIES

The first theorem says that the potential $V(x)$ is only allowed to be a special function.

Theorem 1. *Under the assumption in the introduction, we can conclude*

$$V(x) = \sum_{1 \leq i < j \leq n} u(x_i - x_j) \quad \text{if } W \text{ is of type } A_{n-1},$$

$$V(x) = \sum_{1 \leq i < j \leq n} \left(u(x_i - x_j) + u(x_i + x_j) \right) + \sum_{1 \leq j \leq n} v(x_j) \quad \text{if } W \text{ is of type } B_n,$$

$$V(x) = \sum_{1 \leq i < j \leq n} \left(u(x_i - x_j) + u(x_i + x_j) \right) \quad \text{if } W \text{ is of type } D_n.$$

Here $u(t)$ and $v(t)$ are following functions with complex numbers C_1, C_2, \dots :

If W is of type A_{n-1} with $n \geq 3$,

$$(1) \quad u(t) = C_1 \wp(t) + C_2.$$

If W is of type B_n with $n \geq 3$,

$$(2) \quad \begin{cases} u(t) = C_1 \wp(t) + C_2, \\ v(t) = \frac{C_3 \wp(t)^4 + C_4 \wp(t)^3 + C_5 \wp(t)^2 + C_6 \wp(t) + C_7}{\wp'(t)^2} \end{cases}$$

or

$$(3) \quad u(t) = C_1 t^{-2} + C_2 t^2 + C_3 \quad \text{and} \quad v(t) = C_4 t^{-2} + C_5 t^2 + C_6$$

or

$$(4) \quad u(t) = C_1 \quad \text{and} \quad v(t) \text{ is any even function.}$$

If W is of type D_n with $n \geq 4$, then u is (2) or (3).

If W is of type B_2 , then $(u(t), v(t))$ is (2) or (3) or (4) or

$$(5) \quad \begin{cases} u(t) = \frac{C_3 \wp(\frac{t}{2})^4 + C_4 \wp(\frac{t}{2})^3 + C_5 \wp(\frac{t}{2})^2 + C_6 \wp(\frac{t}{2}) + C_7}{\wp'(\frac{t}{2})^2}, \\ v(t) = C_1 \wp(t) + C_2 \end{cases}$$

or

$$(6) \quad \begin{cases} u(t) = C_1 \wp(t) + C_2 \frac{(\wp(\frac{t}{2}) - e_3)^2}{\wp'(\frac{t}{2})^2} + C_3, \\ v(t) = C_4 \wp(t) + \frac{C_5}{\wp(t) - e_3} + C_6 \end{cases}$$

or

$$(7) \quad v(t) = C_1 \quad \text{and} \quad u(t) \text{ is any even function.}$$

In the above theorem, $\wp(t)$ is the Weierstrass elliptic function $\wp(t|2\omega_1, 2\omega_2)$ with primitive half-periods ω_1 and ω_2 which are allowed to be infinity and e_3 is a complex number satisfying $\wp'^2 = 4(\wp - e_1)(\wp - e_2)(\wp - e_3)$ (cf. [WW]). In particular

$$\wp(t|\sqrt{-1}\pi, \infty) = \sinh^{-2} t + \frac{1}{3} \quad \text{and} \quad \wp(t|\infty, \infty) = t^{-2}.$$

Then we note that $(u(t), v(t))$ in (2) has 9 complex parameters including the periods.

Theorem 2. i) If W is of type B_n , the expression of $V(x)$ by u and v is not unique and then we may assume that the coefficient of $\partial_1 \partial_2$ of P_2 equals $2u(x_1 - x_2) - 2u(x_1 + x_2)$ without changing the commuting algebra $\mathbb{C}[P_1, \dots, P_n]$.

ii) If W is not of type A_{n-1} or if W is of type A_{n-1} and $\text{ord } R_3 < 2$, then $\mathbb{C}[P_1, \dots, P_n]$ is uniquely determined by u or (u, v) .

iii) The commuting differential operators P_1, \dots, P_n exist for P with the potential $V(x)$ defined by u and v of the form (1), (2), (4), (5), (6) and (7) according to the type of W , where C_1, \dots are any complex numbers.

If W is of type A_{n-1} , the commuting differential operators are given by

$$P_k = \sum_{0 \leq j \leq \lfloor \frac{k}{2} \rfloor} \frac{1}{2^j j! (k-2j)!} \sum_{\sigma \in \mathfrak{S}_n} \sigma \left(u(x_1 - x_2) u(x_3 - x_4) \cdots \right. \\ \left. \cdot u(x_{2j-1} - x_{2j}) \partial_{2j+1} \partial_{2j+2} \cdots \partial_k \right)$$

for $k = 1, \dots, n$ (cf. [OP2] and [OS]).

If W is of type B_n and

$$u(t) = C_5 \wp(t), \quad v(t) = \sum_{j=1}^4 C_j \wp(t + \omega_j) - \frac{C_0}{2}$$

with complex numbers C_0, \dots, C_5 and $\omega_3 = -(\omega_1 + \omega_2)$ and $\omega_4 = 0$, then the commuting operators are given by

$$P_n(C_0) = \sum_{k=0}^n \frac{1}{k!(n-k)!} \sum_{\sigma \in \mathfrak{S}_n} \sigma \left(q_{\{1, \dots, k\}} \Delta_{\{k+1, \dots, n\}}^2 \right)$$

(cf. [O]), where

$$\Delta_{\{1, \dots, k\}} = \sum_{0 \leq j \leq \lfloor \frac{k}{2} \rfloor} \frac{1}{2^{k-j} j! (k-2j)!} \sum_{w \in W(B_k)} \varepsilon(w) w \left(u(x_1 - x_2) u(x_3 - x_4) \cdots \right. \\ \left. \cdot u(x_{2j-1} - x_{2j}) \partial_{2j+1} \partial_{2j+2} \cdots \partial_k \right),$$

$$q_{\{1, \dots, k\}} = \sum_{I_1 \amalg \cdots \amalg I_\nu = \{1, \dots, k\}} T_{I_1} \cdots T_{I_\nu}, \quad q_\emptyset = 1,$$

$$T_{\{1, \dots, k\}} = (-C_5)^{k-1} \left(\frac{C_0}{2} T_{\{1, \dots, k\}}^0(1) - \sum_{j=1}^4 C_j T_{\{1, \dots, k\}}^0(\wp(t + \omega_j)) \right),$$

$$T_{\{1, \dots, k\}}^0(\psi) = \sum_{I_1 \amalg \cdots \amalg I_\nu = \{1, \dots, k\}} (-1)^{\nu-1} (\nu-1)! S_{I_1}(\psi) \cdots S_{I_\nu}(\psi),$$

$$S_{\{1, \dots, k\}}(\psi) = \sum_{w \in W(B_k)} w \left(\psi(x_1) \wp(x_1 - x_2) \wp(x_2 - x_3) \cdots \wp(x_{k-1} - x_k) \right).$$

Here $W(B_k)$ and $W(D_k)$ are the Weyl groups of type B_k and D_k , respectively, $W(B_k)$ and $W(D_k)$ and \mathfrak{S}_k are realized as groups of coordinate transformations of \mathbb{R}^k . For $w \in W(B_k)$, $\varepsilon(w) = 1$ if $w \in W(D_k)$ and -1 otherwise, the sums for I_1, \dots, I_ν run over all the partitions of $\{1, \dots, k\}$, and for a subset I of $\{1, \dots, n\}$, we define $\Delta_I = \sigma(\Delta_{\{1, \dots, k\}})$ etc. by $\sigma \in \mathfrak{S}_n$ and $k = \#I$ with $\sigma(\{1, \dots, k\}) = I$.

Expanding $P_n(C_0)$ into a polynomial function of the parameter C_0 , the operators P_j are given by the coefficients of C_0^{n-j} in the expansion. In fact we have $[P_n(C_0), P_n(C'_0)] = 0$.

If W is of type D_n , we have only to put $C_1 = C_2 = C_3 = C_4 = 0$ and $P_n = \Delta_{\{1, \dots, n\}}$ in the above definition. See [OO] for other cases of type B_2 .

Remark 3. i) If (u, v) is of the form (3), P_j do not exist in general and we need operators of higher order (cf. [OP2]).

ii) If (u, v) is given by (4), then $\mathbb{C}[P_1, \dots, P_n]$ equals the totality of \mathfrak{S}_n -invariants in $\mathbb{C}[-\frac{1}{2}\partial_1^2 + v(x_1), \dots, -\frac{1}{2}\partial_n^2 + v(x_n)]$.

iii) If $2\omega_1 = \sqrt{-1}\lambda^{-1}\pi$ and $\omega_2 = \infty$ with $\lambda \neq 0$, (2) is reduced to

$$\begin{cases} u(t) = C'_1 \sinh^{-2} \lambda t + C'_2, \\ v(t) = C'_3 \sinh^{-2} \lambda t + C'_4 \sinh^{-2} 2\lambda t + C'_5 \sinh^2 \lambda t + C'_6 \sinh^2 2\lambda t + C'_7. \end{cases}$$

The commuting differential operators studied by Heckman-Opdam correspond to this case with $C'_5 = C'_6 = 0$. Moreover if $\omega_1 = \omega_2 = \infty$, then (2) is reduced to

$$\begin{cases} u(t) = C'_1 t^{-2} + C'_2, \\ v(t) = C'_3 t^{-2} + C'_4 t^2 + C'_5 t^4 + C'_6 t^6 + C'_7. \end{cases}$$

iv) Some results stated in this note were announced in [Sh]. The precise statements and arguments will be given in [OS], [O] and [OO].

v) Replacing ∂_i , x_j , $[,]$ and ord by $\sqrt{-1}p_i$, q_j , the Poisson bracket $\{ , \}$ and the degree for p , respectively, we have the same statements as in Theorem 1 and Theorem 2, and moreover the operators P_1, \dots, P_n give the integrals of the Hamiltonian corresponding to the Laplacian P (cf. [OP1] for completely integrable classical systems).

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