Phase field models with evaporation & condensation Akira Onuki Kyoto University **Today's talk:**

Part 1) Liquid- Gas : Spreading & Evaporation

Part2) Liquid-solid with inhomogeneous T

Latent heat is crucial in these examples

Phase changes in inhomogeneous T are still unsolved . Usual GL theory is inapplicable.

Part1: Gas-Liquid in one-component fluids: n: number density, e: internal energy density Entropy $S = \int d\mathbf{r} [ns(n,e) - C|\nabla n|^2/2]$ regular (bulk) gradient $\mathcal{E} = \int d\mathbf{r}\hat{e}$ Internal energy internal energy $\hat{e} = e + K |\nabla n|^2 / 2$ density regular (bulk) gradient **Definition of T:** $1/T = \delta S/\delta e$ at fixed n valid for inhomogeneous T

van der Waals
$$(v_0, \epsilon = \text{const})$$

 $s = k_{\mathsf{B}} \ln[(e/n + \epsilon v_0 n)^{d/2} (1/v_0 n - 1)] + const.$

$$e = dnk_{\mathsf{B}}T/2 - \epsilon v_0 n^2$$
$$p = nk_{\mathsf{B}}T/(1 - v_0 n) - \epsilon v_0 n^2$$

Stress tensor for K=0 $\Pi_{ij} = p\delta_{ij} - CT[n\nabla^2 n + (\nabla n)^2/2]\delta_{ij}$ $+ CT\nabla_i n\nabla_j n - \sigma_{ij}$

> reversible stress + viscous stress (η,ζ) Korteweg stress (1901)

$$\frac{\partial}{\partial t}\rho = -\nabla(\rho \mathbf{v}), \ (\rho = mn)$$
$$\rho(\frac{\partial}{\partial t} + \mathbf{v} \cdot \nabla)\mathbf{v} = -\nabla \cdot \Pi - \rho g \mathbf{e}_z$$
$$\frac{\partial}{\partial t} e_T = -\nabla \cdot (e_T + \Pi \cdot)\mathbf{v}$$
$$+\nabla \lambda \nabla T - \rho g v_z, \ (e_T = \hat{e} + \rho v^2/2)$$

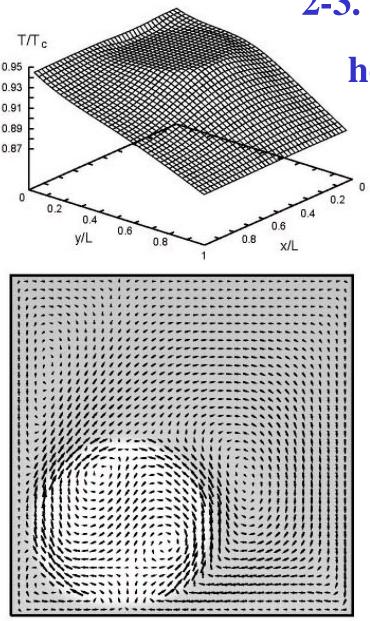
Stress Π contains gradient terms

Entropy production >0 if no heat from outside. We require

 $\frac{\partial}{\partial t} \mathcal{S} = \int d\mathbf{r} \frac{1}{T} [\nabla \cdot \lambda \nabla T + \sum_{ij} \sigma_{ij} \nabla_i v_j]$

Gradient terms are automatically determined

2D simulations (Onuki PRE 2004)



2-3. Bubble is attracted to

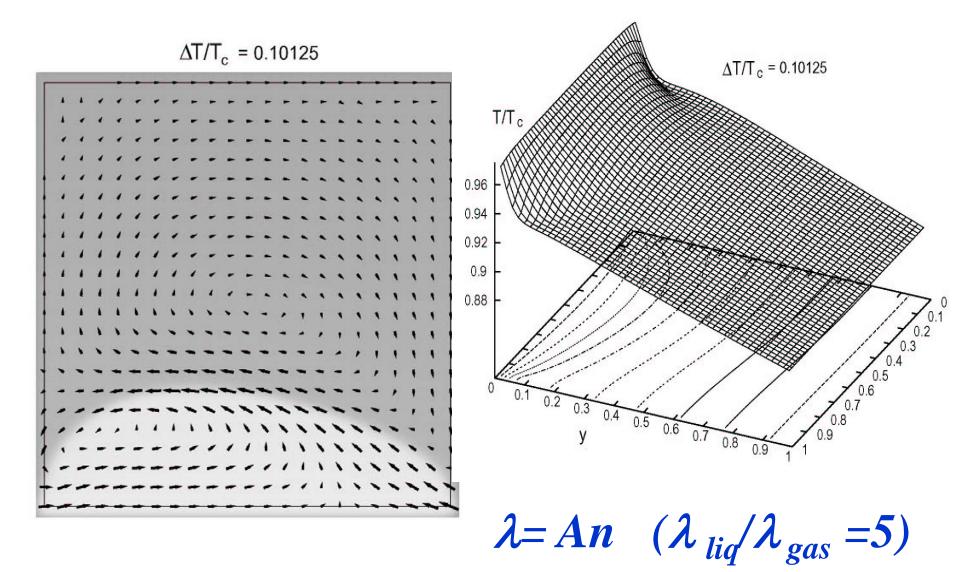
heated wall. Apparent partial

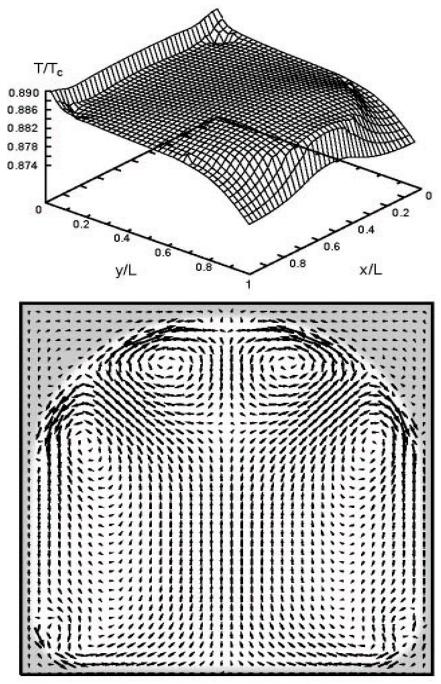
wetting

 $T_{\rm bot}/T_c$ is changed from 0.875 to 0.945 $T_{\rm top}/T_c = 0.875$

Pressure is nearly homogeneous. Interface is nearly on the coexistence curve. T is flat within bubble = Tcx(p) due to latent heat convection

2-4. Wall is covered by gas. Latent heat transport is suppressed Large gradient in gas film





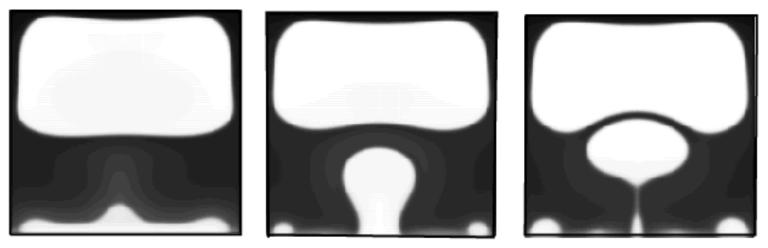
2.5 Efficient heat transport by flow along wetting layer, no gravity

 $\begin{array}{l} \lambda_{\mathrm{liq}} = 5\lambda_{\mathrm{gas}} \\ Nu = \frac{\lambda_{\mathrm{eff}}}{\lambda_{\mathrm{liq}}} = 5 \end{array}$

 $T_{\rm bot}/T_c = 0.895$ $T_{\rm top}/T_c = 0.875$

Steady flow, but drying for larger Tbot. Heat pipe mechanism.

8-1. Boiling: gravity applied, Tt/Tc=0.77, Tb/Tc is changed from 0.94 to 1 at t=34000



t=34600











t=66300

t=71000

t=74100

Simulation: axysymmetric (3D) 600 x 600 Teshigawara &Onuki PRE2011

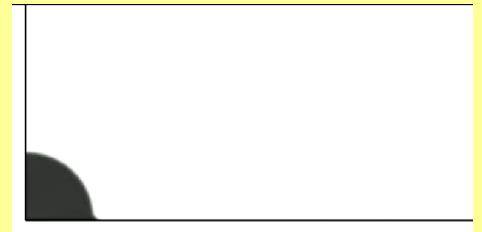
axis ceiling н vapor liquid • temperature z **Bottom of the substrate : given constant Ceiling : given constant** 0 Side wall : insulating • density substrate -Hw **Boundary condition on number density** r substrate in solid side wall $\left. \frac{dn}{dz} \right|_{z=H} = 0$ $C_w \frac{\partial T}{\partial t} = \lambda_w \nabla^2 T,$ **ceiling** $\left. \frac{dn}{dr} \right|_{r} = 0$ thermal conductivity ratio : $\Lambda = \lambda / (nv_0\lambda_w)$

side wall

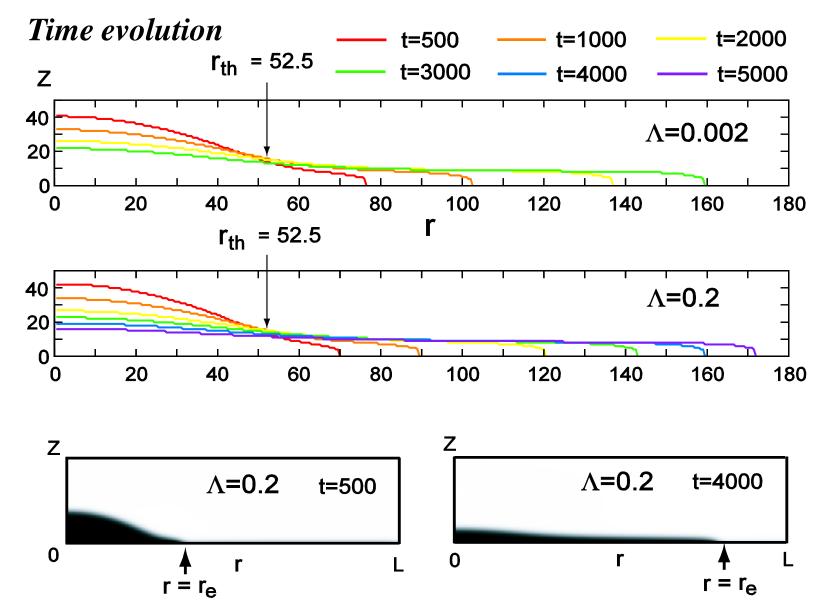
a) Precursor film formation in complete wetting

Plate temperature Tw was not higher than initial liquid temperature

Complete wetting case

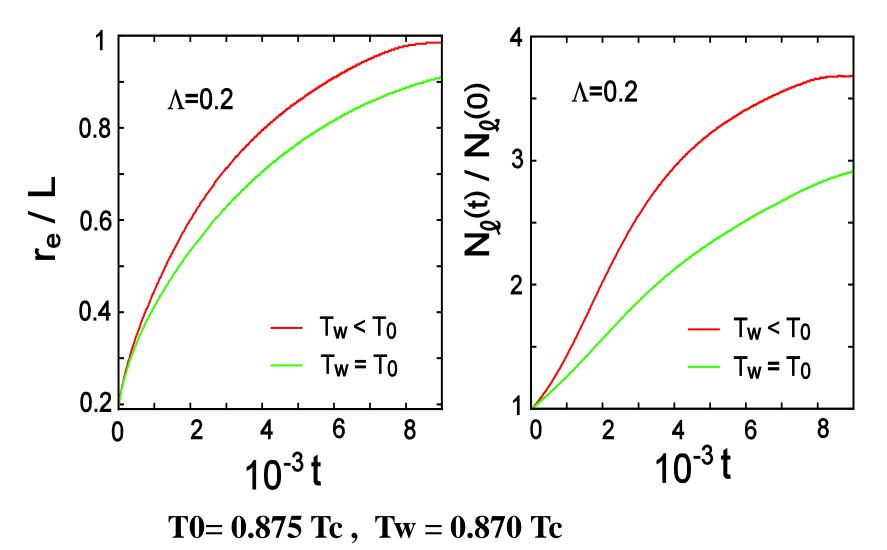


Fast growth of precurserPrecursor filmfilm due to evaporation and condensationLarge wetting parameter $\Phi_1 \Rightarrow$ Precursor film extension

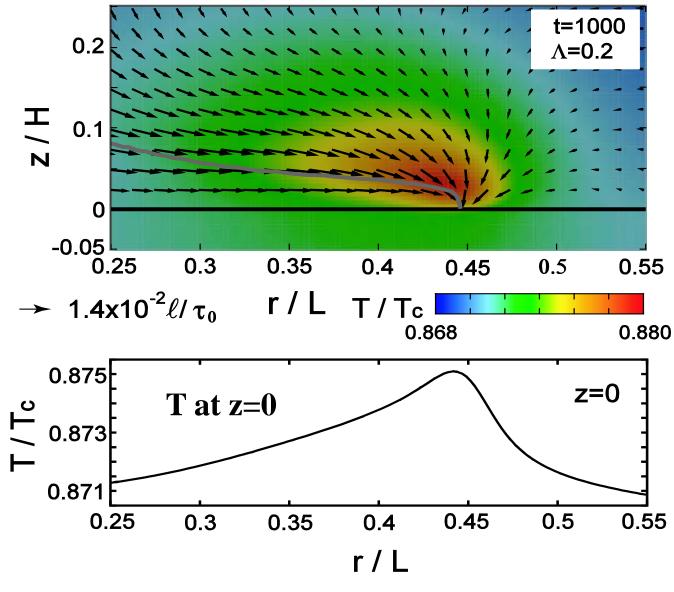


Film thickness = constant. Film expands mainly due to condensation (or flow from the droplet body is smaller).

Condensationoccurs for cooled wallre (t) : edge radiusNi (t) : particle number in liquid

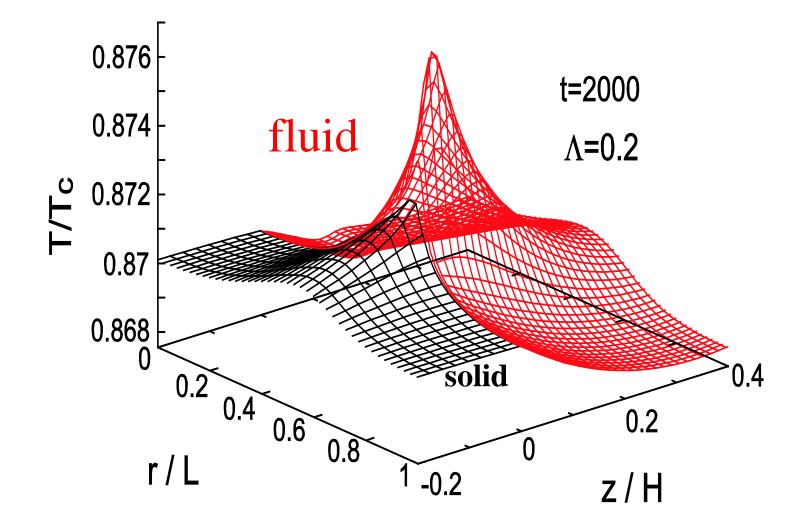


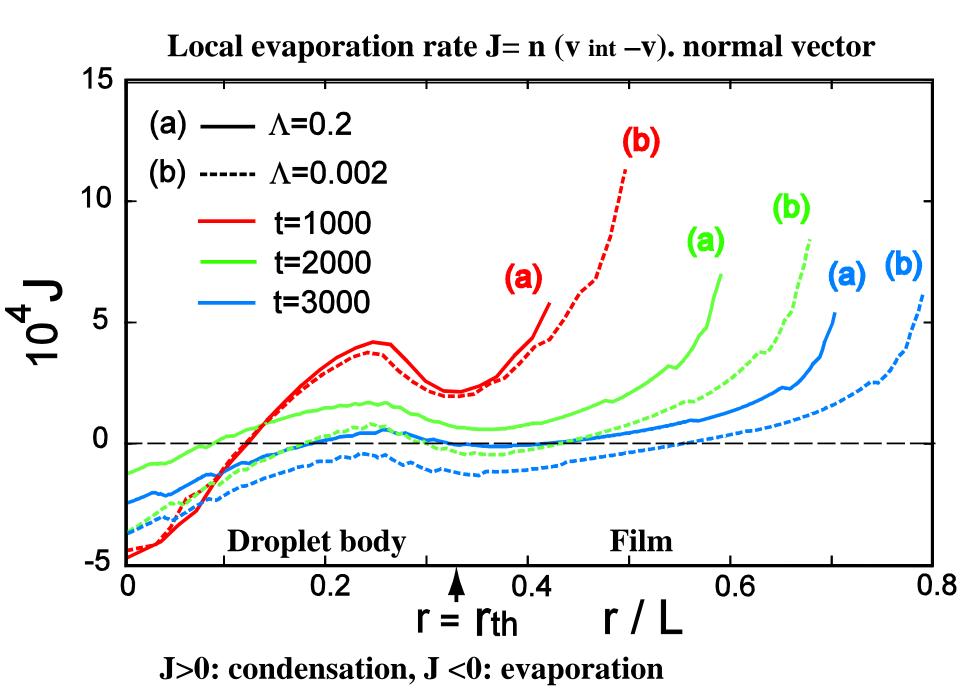
A hot spot appears due to condensation at edge on a cooled wall



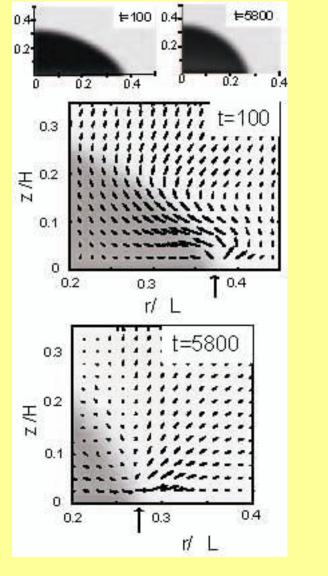
Substrate temperature is inhomogeneous

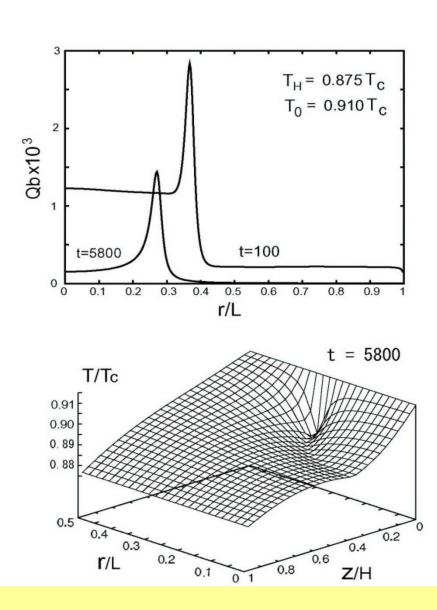
Bird view of temperature around hot spot

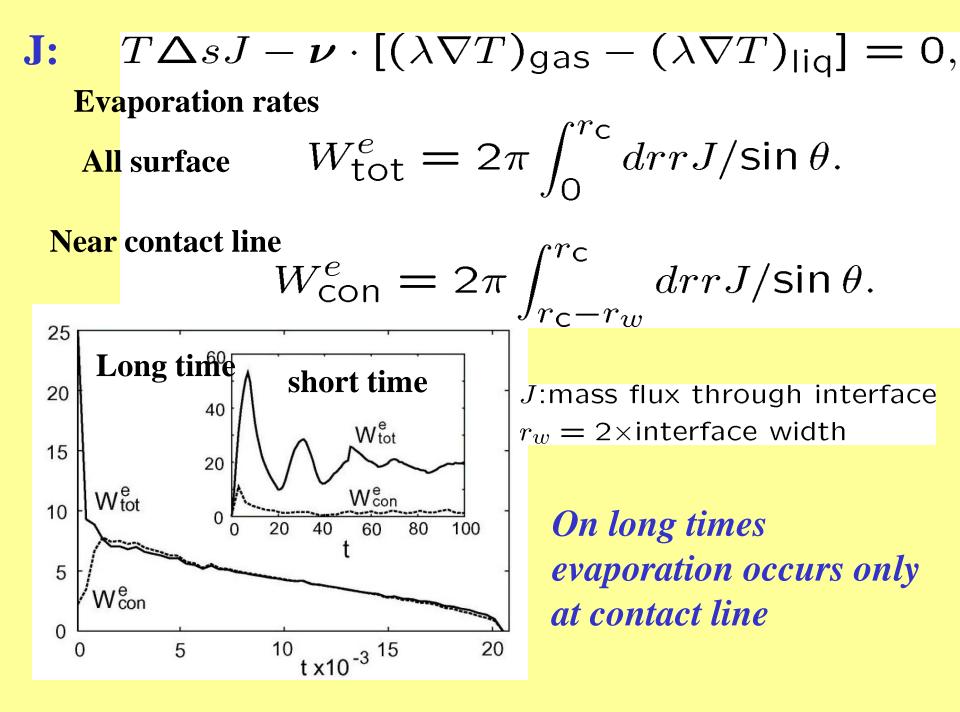




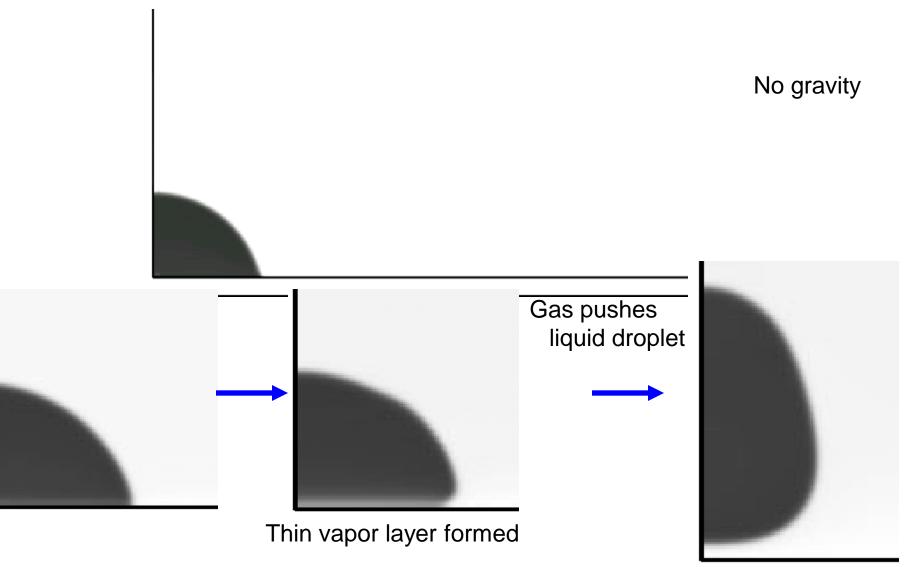
b) Droplet evaporation in partial wetting







Leidenfrost effect: droplet On wall T/Tc is changed **detachment by strong heating** from 0.875 to 0.975



Part 2) solid-liquid (one-component) anisotropic strains e₂ & e₃

Takae & Onuki, Pre (2011)

\phi: Phase field emtropy density : $\hat{S} = S(\rho, e, \phi) - \frac{1}{2}C|\nabla\phi|^2$

energy density : $e_{\rm T} = e + \rho |v|^2 / 2 + G(\phi)(e_2^2 + e_3^2) / 2$

 e_2 :elongational strain, e_3 :shear strain, Here linear elasticity $G = G_0 \phi^2$:shear modulus Expand S up to second order at a reference state:

$$S = S_{0\ell} + \frac{\delta e - \mu_0 \delta \rho}{T_0} - C_V^0 \frac{\tau^2}{2} - \frac{\zeta^2}{2T_0 K_T^0} - W(\phi)$$

We introduce two field variables

$$\tau = (\delta e - \beta_0 \delta \rho) / T_0 C_V^0 - a_1 \theta(\phi)$$

$$\zeta = \delta \rho / \rho_{0\ell} - a_2 \theta(\phi)$$

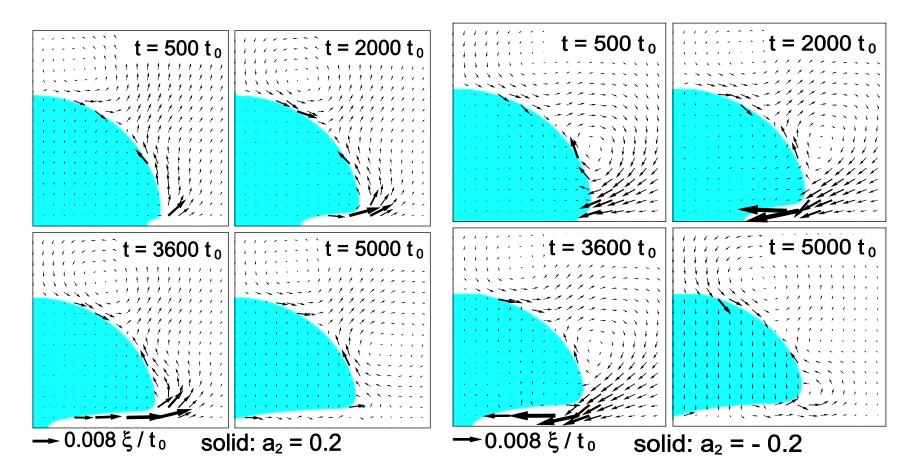
$$W(\phi) = \frac{1}{2}A\phi^2(1-\phi)^2, \quad \theta(\phi) = \phi^2(3-2\phi)$$

Hydrodynamics + elasticity+ phase transition 2D

Phase-field hydrodynamics for solid-liquid:

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ho oldsymbol{v} oldsymbol{v} + \overleftrightarrow{\mathbf{H}} - \overleftrightarrow{\sigma}), \ &rac{\partial e_T}{\partial t} = -
abla \cdot [e_T oldsymbol{v} + (\overleftrightarrow{\mathbf{H}} - \overleftrightarrow{\sigma}) \cdot oldsymbol{v} - \lambda
abla T] \ &rac{\partial }{\partial t} e_2 = -oldsymbol{v} \cdot
abla e_2 +
abla _x v_x -
abla _y v_y \ &rac{\partial }{\partial t} e_3 = -oldsymbol{v} \cdot
abla e_3 +
abla _x v_y +
abla _y v_x. \end{aligned}$$

Example1) "Ice on a hot plate"

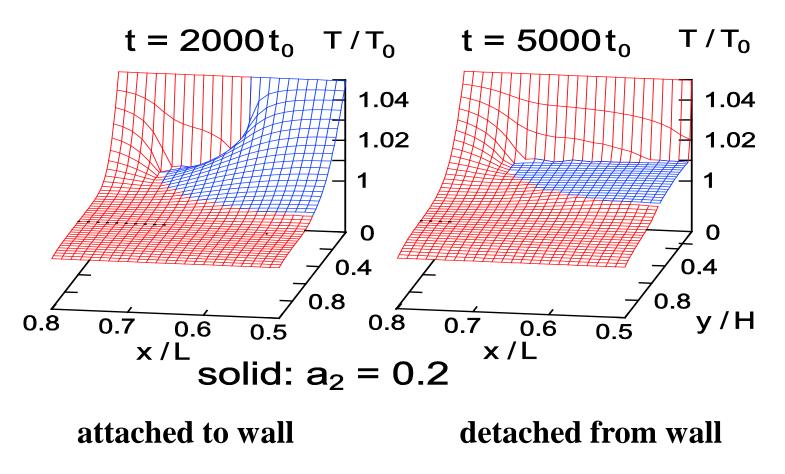


 ρ Isolid > ρ Iiq outgoing flow

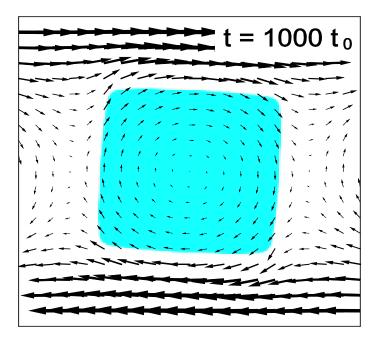
ρliq< ρsolid incoming flow

Temperature around heated ice

Blue: solid, red: liquid

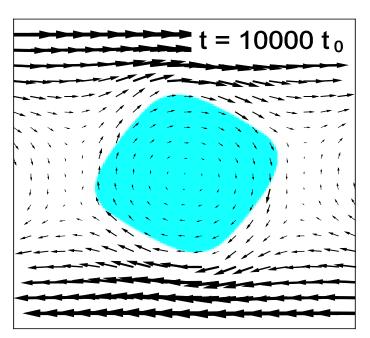


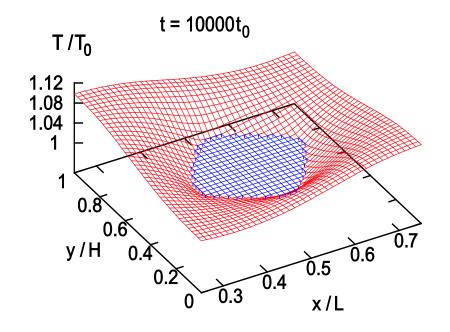
Latent heat cools the liquid region



Example2)Ice cube in warm liquid in shear flow

Latent heat due to melting cools liquid Shear flow accelerates melting





Generalization to account for plasticity Nonlinear elasticity (periodic in e_2 and e_3)

$$e_2 =
abla_x u_x -
abla_y u_y, \quad e_3 =
abla_x u_y +
abla_y u_x.$$

elastic energy = $K(\nabla \cdot \boldsymbol{u})^2/2 + G(\phi)\Phi(e_2,e_3)$

$$\Phi = \frac{1}{6\pi^2} \Big[3 - \cos(2\pi e_2) - \cos(2\pi e_+) - \cos(2\pi e_-) \Big]$$

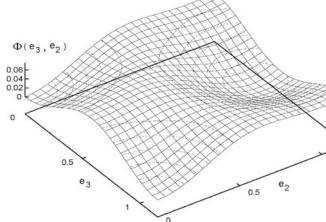
$$(e_{\pm} = (\sqrt{3}e_3 \pm e_2)/2)$$

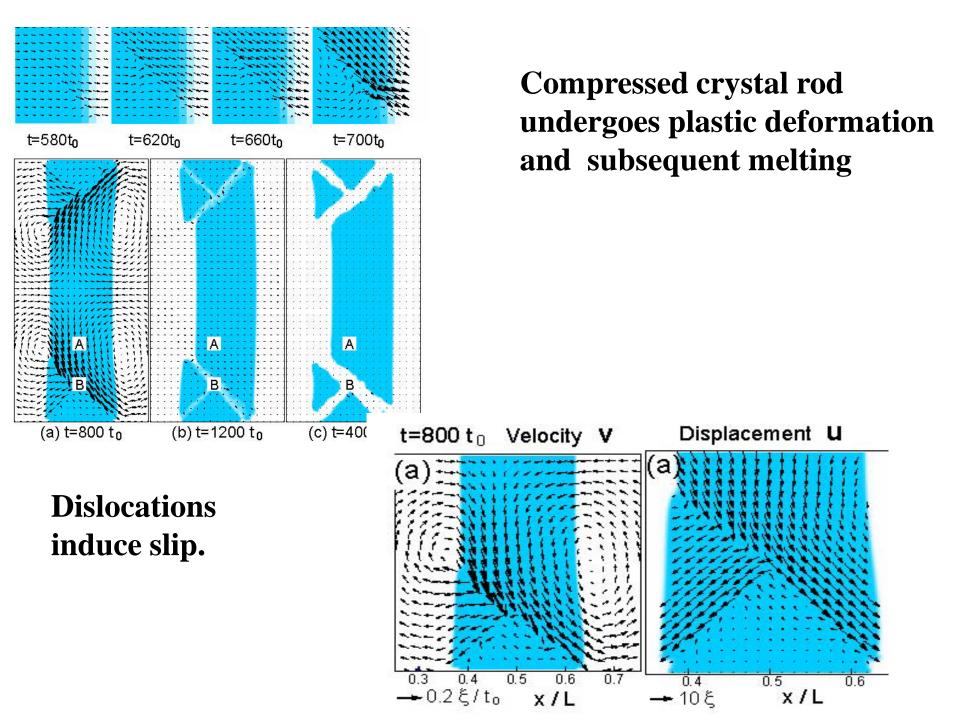
(invariant with respect to $\pm \pi/6$ rotations)

For small e_2 and e_3 , linear elasticit

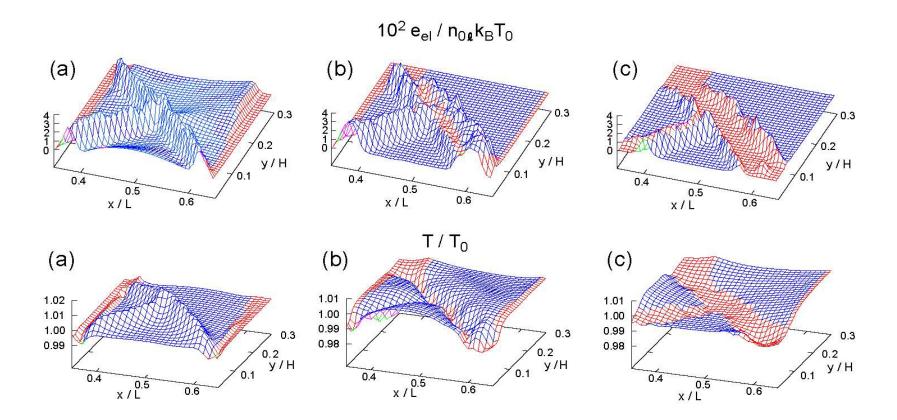
 $\Phi \cong (e_2^2 + e_3^2)/2$

Jumps among multiple minima induce dislocations





Elastic energy density and temperature around plastically deformed region (botom)



Viscous heating

Latent heat cooling upon melting

Summary

1. Two-phase dynamics with evaporation and condensation

2. Wetting dynamics: evaporation, spreading,....

3 Solid-liquid with elasticity in flow

Many important problems in future: In real systems: multi-component fluids ! Boiling (water+air), Marangoni convection