UTMS 2006-32

November 8, 2006

Solution formula for an inverse problem with underdetermining data

by

Yu. E. ANIKONOV and M. YAMAMOTO



UNIVERSITY OF TOKYO

GRADUATE SCHOOL OF MATHEMATICAL SCIENCES KOMABA, TOKYO, JAPAN

¹ Yu. E. Anikonov and ² M.Yamamoto SOLUTION FORMULA FOR AN INVERSE PROBLEM WITH UNDERDETERMINING DATA

¹ Sobolev Institute of Mathematics, Acad. Koptyug prospekt 4 630090 Novosibirsk Russia anikon@math.nsc.ru
² Department of Mathematical Sciences, The University of Tokyo 3-8-1 Komaba, Meguro, Tokyo 153 Japan myama@ms.u-tokyo.ac.jp tel: +81-3-5465-8328 fax: +81-3-5465-7011

§1. Introduction and Key Lemma.

In an inverse problem, we are required to determine coefficients in a partial differential equation in order that the solution to the differential equation realizes prescribed data. As the mathematical topics for an inverse problem, we mention the uniqueness and the stability, and additionally the existence of a solution to the inverse problem is important. Usually the solution to an inverse problem is given not by formulae which are involved by algebraic operations and calculi, but is found through limit processes such as iterations (for example, as a solution to an operator equation of the second kind). In this paper, we will show a formula for solutions to an inverse problem which is attached with underdetermining data.

Our formulation for the inverse problem is underdeterming and so cannot guarantee the uniqueness for solutions to the inverse problem. Hence our formula gives "one" solution to the inverse problem under consideration, and does not describe all possible solutions but includes sufficiently many solutions in the sense that it admits a family of coefficients parameterized by free functions in the spatial variable.

Among various inverse problems, for our approach, we will mainly discuss an inverse problem with data at final time. For example, in Bouchouev and Isakov [3], Isakov [5], such an inverse problem for the Black - Scholes equation is considered: Determine a, b, c by $w(x, T), x \in I$, in

$$\alpha \frac{\partial w}{\partial t} = a(x) \frac{\partial^2 w}{\partial x^2} + b(x) \frac{\partial w}{\partial x} + c(x)w.$$

Here T > 0 is fixed, I is an interval and $\alpha = const > 0$. As for inverse problems with data at final time, we can further refer to Choulli and Yamamoto [4], Prilepko, Orlovsky and Vasin [6] and the references therein.

In the present paper, we construct a family of solutions of the inverse problems, which is based on representation of solutions and coefficients.

Let $D \in \mathbb{R}^n$ be a domain and let us consider an evolution equation of the following type:

$$\sum_{k=1}^{m} \alpha_k(v(x)) \frac{\partial^k w}{\partial t^k} = \sum_{i,j=1}^{n} a_{ij}(x) \frac{\partial^2 w}{\partial x_i \partial x_j} + \sum_{j=1}^{n} b_j(x) \frac{\partial w}{\partial x_j} + c(x)w, \quad (1.1)$$

where $x = (x_1, ..., x_n) \in D \subset \mathbb{R}^n$, $0 \leq t \leq T$, $\alpha_k : \mathbb{R}^n \longrightarrow \mathbb{R}$ are smooth functions for $1 \leq k \leq m$, $v(x) = (v_1(x), \ldots v_n(x))$ is a differentiable vector-valued function such that

$$\left|\frac{\partial(v_1,...,v_n)}{\partial(x_1,...,x_n)}\right| \neq 0, x \in \overline{D}$$

and $a_{ij} = a_{ji}, 1 \leq i, j \leq n$.

First we formulate a general approach for obtaining a representation formula of solution w(x,t) and coefficients $a_{ij}(x)$, $b_j(x)$, c(x).

Note that in the case where

$$\alpha_1 = \alpha = const > 0, \quad \alpha_k = 0, \quad k = 2, 3, \dots, m,$$

and

$$\sum_{i,j=1}^{n} a_{ij}\eta_j\eta_j > 0, \quad \eta \in \mathbb{R}^n, \eta \neq 0,$$

equation (1.1) becomes parabolic:

$$\alpha \frac{\partial w}{\partial t} = \sum_{i,j=1}^{n} a_{ij}(x) \frac{\partial^2 w}{\partial x_i \partial x_j} + \sum_{j=1}^{n} b_j(x) \frac{\partial w}{\partial x_j} + c(x)w.$$

In the case of

$$\alpha_2 = \alpha = const, \quad \alpha_k = 0, k = 1, 3, \dots, m$$

equation (1.1) becomes hyperbolic when $\alpha > 0$, and elliptic when $\alpha < 0$. Moreover if $\alpha_1 = -\sqrt{-1}$ and $\alpha_2 = \cdots = \alpha_m = 0$, then (1.1) is the Schrödinger equation.

Lemma 1. Let $D_1 \subset \mathbb{R}^n$ be a domain and $\beta_{k\ell}$, β_k , $\beta \in C^1(D_1)$, $1 \leq k, \ell \leq n$ such that $\beta_{k\ell} = \beta_{\ell k}$, be given, and $u \in C^2(D)$ be a given function such that $u(x) \neq 0$ for any $x \in \overline{D}$. Let

$$F(y,t), \quad y \in D_1 \subset \mathbb{R}^n, 0 \le t \le T$$

satisfy the equation

$$\sum_{k=1}^{m} \alpha_k(y) \frac{\partial^k F}{\partial t^k} = \sum_{k,\ell=1}^{n} \beta_{k\ell}(y) \frac{\partial^2 F}{\partial y_\ell \partial y_k} + \sum_{k=1}^{n} \beta_k(y) \frac{\partial F}{\partial y_k} + \beta(y)F.$$
(1.2)

Then the function

$$w(x,t) = u(x)F(v(x),t)$$

and the coefficients $a_{ij}(x)$, $b_j(x)$, c(x) consecutively defined by linear systems of algebraic equations (1.3), (1.4), (1.5), satisfy equation (1.1):

$$\sum_{i,j=1}^{n} a_{ij}(x) \frac{\partial v_k}{\partial x_i} \frac{\partial v_\ell}{\partial x_j} = \beta_{k\ell}(v(x))$$
(1.3)

$$\sum_{j=1}^{n} b_j(x) \frac{\partial v_k}{\partial x_j} = \frac{1}{u} \left[u\beta_k(v) - \sum_{i,j=1}^{n} a_{ij} \left(\frac{\partial v_k}{\partial x_i} \frac{\partial u}{\partial x_j} + \frac{\partial v_k}{\partial x_j} \frac{\partial u}{\partial x_i} + u \frac{\partial^2 v_k}{\partial x_i \partial x_j} \right) \right],$$

$$(1.4)$$

$$c(x) = \frac{1}{u} \left[u\beta(v) - \left(\sum_{i=1}^{n} a_{ii} \frac{\partial^2 u}{\partial x_i} + \sum_{i=1}^{n} b_i(x) \frac{\partial u}{\partial x_i} \right) \right].$$

$$(1.5)$$

$$c(x) = \frac{1}{u} \left[u\beta(v) - \left(\sum_{i,j=1}^{n} a_{ij} \frac{\partial^2 u}{\partial x_i \partial x_j} + \sum_{j=1}^{n} b_j(x) \frac{\partial u}{\partial x_j} \right) \right].$$
(1.5)

The proof of the lemma is done by direct substitution of the solution w(x,t) = u(x)F(v(x),t) into equation (1.1) in terms of (1.2) - (1.5).

We note that if we choose v(x) = x, u(x) = 1, $\beta_{k\ell} = a_{k\ell}$, $\beta_k = a_k$ and $\beta = c$ for $1 \le k, \ell \le n$, then (1.3) - (1.5) are true.

After choosing v = v(x), we set $\mathcal{A} = \{u, \{\beta_{k\ell}, \beta_k\}_{1 \le k, \ell \le n}, \beta\}$. Then, by $a_{ij}(\mathcal{A}), b_j(\mathcal{A}), c(\mathcal{A}), 1 \le i, j \le n$, we denote a_{ij}, b_j and c defined by (1.3) - (1.5). We note that \mathcal{A} is composed of $\frac{n^2+3n+4}{2}$ functions in x. Then, by Lemma 1 we can represent $\frac{n^2+3n+2}{2}$ functions a_{ij}, b_j, c in x by $\frac{n^2+3n+4}{2}$ functions in x. In other words, our lemma gives representation formulae of coefficients a_{ij}, b_j, c which contains $1\left(=\frac{n^2+3n+4}{2}-\frac{n^2+3n+2}{2}\right)$ free function in x. Thus by our representation formula, we can give a pair of solution (a_{ij}, b_j, c) which realizes one extra data $w(x, T) = w_1(x), x \in D$. We notice that our formula (1.3) - (1.5) are not involved with limit processes.

In particular, if coefficients α_k , β_{kl} , β_k , β are constant in (1.2), then in some cases it is possible to represent F(y, t) which is a solution to an initial value problem. Moreover, if we a priori know F(y, t), then problem of search for coefficients and solution obviously turns to the determination of only functions u(x), $v(x) = (v_1(x), \ldots, v_n(x))$.

We will consider a one-dimensional variant of Lemma 1 more precisely. Let w(x,t) satisfy the following equation:

$$\sum_{k=1}^{m} \alpha_k(v(x)) \frac{\partial^k w}{\partial t^k} = a(x) \frac{\partial^2 w}{\partial x^2} + b(x) \frac{\partial w}{\partial x} + c(x)w, \qquad (1.6)$$

 $x_0 \le x \le x_1, \quad 0 < t < T,$

where $\alpha_k(y), y \in \mathbb{R}^1, v(x), a(x), b(x), c(x)$ are some differentiable functions and α_k may be constant.

Lemma 2. Let

$$u(x), v(x), \alpha_k(y), \beta_1(y), \beta_2(y), \beta_3(y), k = 1, 2, \dots, m$$

be some twice differentiable functions and let

$$x_0 \le x \le x_1, \quad y \in \mathbb{R}^1, \quad u(x) \ne 0, \quad \frac{dv}{dx} \equiv v'(x) \ne 0.$$

We assume that $F(y, t), y \in \mathbb{R}^1, 0 < t < T$, satisfies the equation

$$\sum_{k=1}^{m} \alpha_k(y) \frac{\partial^k F}{\partial t^k} = \beta_1(y) \frac{\partial^2 F}{\partial y^2} + \beta_2(y) \frac{\partial F}{\partial y} + \beta_3(y) F, \quad y \in \mathbb{R}, \, t > 0.$$
(1.7)

Then the functions w(x,t), a(x), b(x), c(x) defined by the following formulae

$$w(x,t) = u(x)F(v(x),t), \quad a(x) = \frac{\beta_1(v(x))}{v'^2},$$
$$b(x) = \frac{v'^2 u \beta_2(v) - \beta_1(v)(2u'v' + uv'')}{uv'^3}$$
$$c(x) = \frac{u^2 v'^3 \beta_3(v) - uu''v'\beta_1(v) - u'v'^2 u \beta_2(v) + 2u'^2 v'\beta_1 + u'uv''\beta_1}{u^2 v'^3}$$

satisfy the equation

$$\sum_{k=1}^{m} \alpha_k(u(x)) \frac{\partial^k w}{\partial t^k} = a(x) \frac{\partial^2 w}{\partial x^2} + b(x) \frac{\partial w}{\partial x} + c(x) w.$$

As for other approaches to inverse problems by means of formulae, we refer to Anikonov [1], [2].

§2. One-dimensional parabolic inverse problem with data at final time.

As an example of using this way we consider an inverse problem for a onedimensional parabolic equation or a modifed Black-Scholes equation (e.g., [3]): Find functions

$$w(x,t), a(x), b(x), c(x), 0 \le t \le T, x \in \mathbb{R},$$

such that

$$\alpha \frac{\partial w}{\partial t} = a(x) \frac{\partial^2 w}{\partial x^2} + b(x) \frac{\partial w}{\partial x} + c(x)w, \quad x \in \mathbb{R}, \, t > 0$$
(2.1)

and

$$w\Big|_{t=0} = w_0(x), w\Big|_{t=T} = w_1(x), \quad x \in \mathbb{R}.$$
 (2.2)

More precisely, represent the four functions w(x,t), a(x), b(x), c(x) by $w_0(x)$, $w_1(x)$ and one real-valued auxiliary function. Here we assume that

$$w_0(x) > 0, \qquad x \in \mathbb{R}. \tag{2.3}$$

Then we note that $w_1(x) > 0$ for $x \in \mathbb{R}$ by the maximum principle.

Theorem. We choose $\delta \in \mathbb{R}$ and a smooth function $f_0 > 0$ for $x \in \mathbb{R}$ such that

$$\Phi(x) = \frac{f_0(x)}{\int_{-\infty}^{\infty} \frac{1}{2\sqrt{\pi T}} \exp\left(-\frac{y^2}{4T}\right) f_0(x-y) dy}$$
(2.4)

is an injective function in $x \ge \delta$. Furthermore we assume that

$$\left[0, \sup_{x \in \mathbb{R}} \frac{w_0(x)}{w_1(x)}\right] \subset \Phi([\delta, \infty)).$$
(2.5)

We set

$$v(x) = \Phi^{-1}\left(\frac{w_0(x)}{w_1(x)}\right), \quad u(x) = \frac{w_1(x)}{w_1(v(x))}, \quad x \in \mathbb{R}.$$
 (2.6)

Then

$$a(x) = \frac{1}{(v'(x))^2},$$

$$b(x) = -\frac{2u'(x)v'(x) + u(x)v''(x)}{u(x)(v'(x))^3},$$

$$c(x) = \frac{-u(x)v'(x)u''(x) + 2(u'(x))^2v'(x) + u(x)u'(x)v''(x)}{(u(x))^2(v'(x))^3},$$

$$w(x,t) = \frac{u(x)}{2\sqrt{\pi t}} \int_{-\infty}^{\infty} \exp\left(-\frac{(x-y)^2}{4t}\right) w_0(v(y))dy, \quad x \in \mathbb{R}, t > 0,$$

$$c(x) = \frac{u(x)}{2\sqrt{\pi t}} \int_{-\infty}^{\infty} \exp\left(-\frac{(x-y)^2}{4t}\right) w_0(v(y))dy, \quad x \in \mathbb{R}, t > 0,$$

satisfies (2.1) and (2.2).

Example. Setting

$$f_0(x) = e^{-x^2}, \qquad x > 0,$$

we see that for any $\delta > 0$, the function Φ is injective in $x \in \mathbb{R}$. In fact,

$$\int_{-\infty}^{\infty} \frac{1}{2\sqrt{\pi T}} \exp\left(-\frac{y^2}{4T}\right) e^{-(x-y)^2} dy = \frac{1}{\sqrt{4T+1}} \exp\left(-\frac{x^2}{4T+1}\right),$$

and so

$$\Phi(x) = \sqrt{4T+1} \exp\left(-\frac{4Tx^2}{4T+1}\right).$$

Thus under the assumption that

$$\sup_{x\in\mathbb{R}}\frac{w_0(x)}{w_1(x)}<\sqrt{4T+1},$$

we can rewrite (2.6) as

$$v(x) = \left\{\frac{4T+1}{4T}\log\left(\sqrt{4T+1}\frac{w_1(x)}{w_0(x)}\right)\right\}^{\frac{1}{2}}$$

and

$$u(x) = \frac{w_1(x)}{w_1(v(x))},$$

so that the conclusion of the theorem holds.

Proof of Theorem. In Lemma 2, we set $\alpha_1 = 1$, $\alpha_2 = \cdots = \alpha_m = 0$, $\beta_1 = 1$, $\beta_2 = \beta_3 = 0$. Then

$$F(y,t) = \frac{1}{2\sqrt{\pi t}} \int_{-\infty}^{\infty} e^{-\frac{y^2}{4t}} f_0(x-y) dy$$

satisfies (1.7). We have

$$w_0(x) = u(x)f_0(v(x)), \quad w_1(x) = u(x)F(v(x),T).$$

Eliminating u(x) in these equations, we obtain

$$\frac{w_0(x)}{w_1(x)} = \frac{F(v(x), 0)}{F(v(x), T)} = \Phi(v(x)).$$

Therefore we have

$$v(x) = \Phi^{-1}\left(\frac{w_0(x)}{w_1(x)}\right)$$

and

$$u(x) = \frac{w_0(x)}{F\left(\Phi^{-1}\left(\frac{w_0(x)}{w_1(x)}\right), 0\right)} = \frac{w_1(x)}{F\left(\Phi^{-1}\left(\frac{w_0(x)}{w_1(x)}\right), 0\right)}$$

Hence, by Lemma 2, the proof of the theorem is complete.

ACKNOWLEDGEMENTS.

This work has been done during the stay of the first author at Graduate School of Mathematical Sciences of the University of Tokyo February of 2004 which was supported by the 21st Center of Excellence Program and he gratefully acknowledges the support. The second author was partly supported by Grants 15340027 from the Japan Society for the Promotion of Science and Grant 17654019 from the Ministry of Education, Cultures, Sports and Technology.

REFERENCES

1. Yu. E. Anikonov, Formulas in Inverse and Ill-posed Problems, VSP, Utrech, 1997.

2. Yu. E. Anikonov, Inverse Problems for Kinetic and Other Evolution Equations, VSP, Utrech, 2001.

3. I. Bouchouev and V.Isakov, Uniqueness, stability, and numerical methods for the the inverse problem that arises in financial markets, Inverse Problems, **15**(1999), R95-R116.

4. M. Choulli and M. Yamamoto, Generic well-posedness of an inverse parabolic problem - the Hölder space approach, Inverse Problems, **12** (1996), 195–205.

5. V.Isakov, The inverse problem of option pricing, in Recent Development in Theories and Numerics, World Scientific, Singapore, 2003, pp.47-55.

6. A.I. Prilepko, D.G. Orlovsky and I.A. Vasin, Methods for Solving Inverse Problems in Mathematical Physics, Maecal Dekker, New York, 2000. Preprint Series, Graduate School of Mathematical Sciences, The University of Tokyo

UTMS

- 2006–21 Yoshihiro Sawano and hitoshi Tanaka : A quarkonial decomposition of Besov-Morrey spaces and Triebel-Lizorkin-Morrey spaces.
- 2006–22 Victor Isakov, Jenn-Nan Wang and Masahiro Yamamoto: An inverse problem for a dynamical Lamé system with residual stress.
- 2006–23 Oleg Yu. Imanuvilov, Victor Isakov and Masahiro Yamamoto: New realization of the pseudoconvexity and its application to an inverse problem.
- 2006–24 Kazuki Hiroe and Takayuki Oda: Hecke-Siegel's pull back formula for the Epstein zeta function with a harmonic polynomial.
- 2006–25 Takefumi Igarashi and Noriaki Umeda: Existence and nonexistence of global solutions in time for a reaction-diffusion system with inhomogeneous terms.
- 2006–26 Shigeo Kusuoka: A remark on law invariant convex risk measures.
- 2006–27 Ganghua Yuan and Masahiro Yamamoto: Lipschitz stability in the determination of principal parts of a parabolic equation by boundary observations.
- 2006–28 Akishi Kato: Zonotopes and four-dimensional superconformal field theories.
- 2006–29 Yasufumi Osajima: The asymptotic expansion formula of implied volatility for dynamic SABR model and FX hybrid model.
- 2006–30 Noriaki Umeda: On existence and nonexistence global solutions of reactiondiffusion equations.
- 2006–31 Takayuki Oda and Masao Tsuzuki: The secondary spherical functions and automorphic green currents for certain symmetric pairs.
- 2006–32 Yu. E. Anikonov and M. Yamamoto: Solution formula for an inverse problem with underdetermining data.

The Graduate School of Mathematical Sciences was established in the University of Tokyo in April, 1992. Formerly there were two departments of mathematics in the University of Tokyo: one in the Faculty of Science and the other in the College of Arts and Sciences. All faculty members of these two departments have moved to the new graduate school, as well as several members of the Department of Pure and Applied Sciences in the College of Arts and Sciences. In January, 1993, the preprint series of the former two departments of mathematics were unified as the Preprint Series of the Graduate School of Mathematical Sciences, The University of Tokyo. For the information about the preprint series, please write to the preprint series office.

ADDRESS:

Graduate School of Mathematical Sciences, The University of Tokyo 3–8–1 Komaba Meguro-ku, Tokyo 153-8914, JAPAN TEL +81-3-5465-7001 FAX +81-3-5465-7012